

Photoresponse of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ nanostrip

Hiroyuki Shibata¹, Naoto Kirigane¹, Kentaro Fukao¹, Daisuke Sakai¹, Shinichi Karimoto², and Hideki Yamamoto²

¹ Kitami Institute of Technology, 165 Koen-cho, Kitami, Hokkaido 090-8507, Japan

² NTT Basic Research Laboratories, 3-1, Wakamiya, Morinosato, Atsugi-shi, Kanagawa 243-0198, Japan

E-mail: shibathr@mail.kitami-it.ac.jp

Abstract

We report the fabrication and photoresponse of 5nm thick $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ nanostrip with a width of 100nm. The I-V characteristics of the nanostrip show a hysteresis and a sharp voltage jump at I_c . The $J_c(3\text{K})$ of the nanostrip is $2.3 \times 10^7 \text{ A/cm}^2$. The nanostrip exhibits photoresponse signals when illuminated by a pulse laser at 1560 nm wavelength with a bias current just below I_c . The height of the signal reduces as the optical intensity decreases and disappears below -10 dBm. The signal also decreases as the temperature increases, but it exists up to 30 K. These results suggest the possibility of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ nanostrip as single-photon detector (SSPD, SNSPD) working at high temperature by further reducing the cross-section of the strip.

1. Introduction

High-performance single-photon detector in the infrared region is required in many fields such as quantum information, quantum optics, space-to-ground communications, and bio sensing [1]. For this purpose, superconducting single-photon detector (SSPDs, SNSPDs) has been attracted much attention in the last decade [1,2]. Now the SSPDs with the detection efficiency close to unity (93%) have been reported by several groups [3,4]. It is also noted that the dark count rate of SSPD has been reached to a theoretical limit of background blackbody radiation [5,6]. By introducing cold bandpass filters with 100 GHz bandwidth, the dark count rate of SSPD decreases to 10^{-4} cps [6]. The SSPD with ultralow dark count rate is particularly important for the long distance quantum key distribution (QKD) over 300 km [7].

Although the performance of SSPDs overcomes to the other single-photon detectors such as InGaAs avalanche photodiode, a drawback is a low operating temperature. The operating temperature is below 3 K for NbN based SSPD, and below 1 K for WSi based SSPD, depending on the T_c of the materials [1-6]. To increase the operating temperature, we have to use higher T_c materials for SSPDs. Recently, we have developed the SSPDs based on MgB₂ ($T_c = 39$ K) [8], and it has been demonstrated that the MgB₂-SSPD exhibits the capacity of detecting single photons up to 10 K [9,10]. To further increase the operating temperature, it is inevitable to use cuprate superconductors for SSPDs.

One of the key technological issues for the realization of cuprate based SSPDs is the ultrathin film growth. We need high quality ultrathin films for SSPDs; with the thickness of about 5nm, a sharp superconducting transition, and a smooth surface. There have been several reports towards the realization of SSPDs using YBa₂Cu₃O_{7- δ} films with T_c about 90 K [11-15]. However, as the thickness decreases to 5 nm, the quality of the films strongly deteriorates [14]. Arpaia et al. fabricated the YBa₂Cu₃O_{7- δ} parallel type meander structure down to 100 nm wide \times 50 nm thick and observed the photoresponse of the meanders [13]. Amari et al. fabricated the YBa₂Cu₃O_{7- δ} meanders with 100 nm wide \times 30 nm thick by ion irradiation [15]. It is necessary to further reduce the thickness of the film for realizing the cuprate based SSPDs.

Among the cuprates, there is a La_{1.85}Sr_{0.15}CuO₄ with T_c of 37 K. It has been reported that it is possible to grow high quality epitaxial films of La_{1.85}Sr_{0.15}CuO₄ with the thickness of 5 nm on a LaSrAlO₄ substrate [16-18]. So, if we use La_{1.85}Sr_{0.15}CuO₄ instead of YBa₂Cu₃O_{7- δ} , we can solve the problem of ultrathin film growth for SSPDs.

Here, we report the fabrication of La_{1.85}Sr_{0.15}CuO₄ nanostrip with the size of 100 nm wide \times 10 μ m long \times 5 nm thick. From the transport measurement, it is revealed that the strip has good superconducting characteristics. The strip shows a photoresponse up to 30 K, suggesting the possible single-photon detection at high temperature in the future.

2. Fabrication

Ultrathin films of La_{1.85}Sr_{0.15}CuO₄ were synthesized in a custom-designed ultrahigh-vacuum chamber (base pressure $<10^{-7}$ Pa) from metal sources by means of reactive co-evaporation [16,17,19].

For oxidation of the film, RF activated atomic oxygen (input power of 300 W and oxygen flow rate of 1 sccm) was used. During the growth, the evaporation beam flux of each element was controlled by electron impact emission spectroscopy (EIES) and the film was monitored by reflection high-energy electron diffraction (RHEED). The growth rate was 9 nm/min and the growth temperature is 670 °C. The 5 nm thick thin films with (001) orientation were epitaxially grown on (001) LaSrAlO₄ substrates. Details of our growth technique have been published elsewhere [16,17,19].

Figure 1 shows the resistivity vs temperature (ρ -T) curves of La_{1.85}Sr_{0.15}CuO₄ thin film with the thickness of 5 nm. As shown in the inset of Fig. 1, T_c ($R = 0$) is 41.6 K, which is higher than the value (37 K) of bulk specimens. This is due to the compressive strain effect caused by the lattice mismatch between film and substrate [16,17]. The in-plane lattice parameter of LaSrAlO₄ is 0.3756 nm, which is about 5% shorter than the 0.3777 nm of La_{1.85}Sr_{0.15}CuO₄. Now, we can use very high quality ultrathin films with T_c above 40 K for the development of SSPDs.

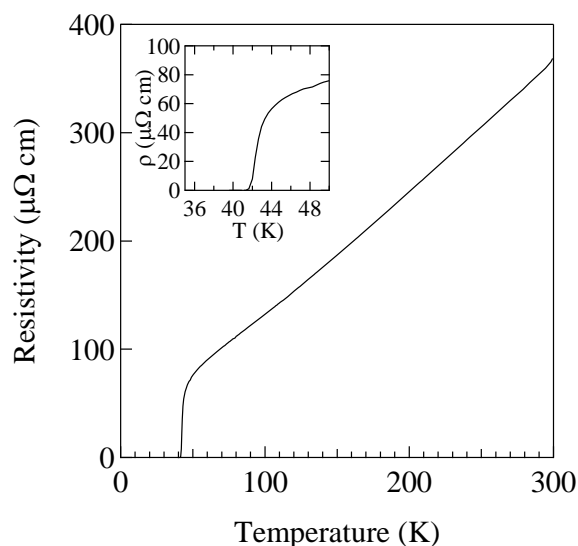


Figure 1. Temperature dependence of the resistivity of 5 nm thick La_{1.85}Sr_{0.15}CuO₄ thin films. Insets: superconducting transition.

For the nanofabrication of cuprates, Au passivation is commonly used to protect the damage. However, it is impossible to use conductive materials as a passivation layer in the case of SSPDs. Here, we use AlN for the passivation of La_{1.85}Sr_{0.15}CuO₄ films. The AlN film has been known to be useful as a passivation layer of MgB₂-SSPDs [20]. The single line pattern with 10 μm long and 100 nm wide was fabricated on silicon-based negative resist (SNR) by e-beam lithography. Figure 2 shows scanning electron microscopy (SEM) images of the pattern. Then the pattern was transferred to the La_{1.85}Sr_{0.15}CuO₄ film using standard Ar ion milling. In this way, we can fabricate La_{1.85}Sr_{0.15}CuO₄ strip with the size of 100 nm wide × 10 μm long × 5 nm thick.

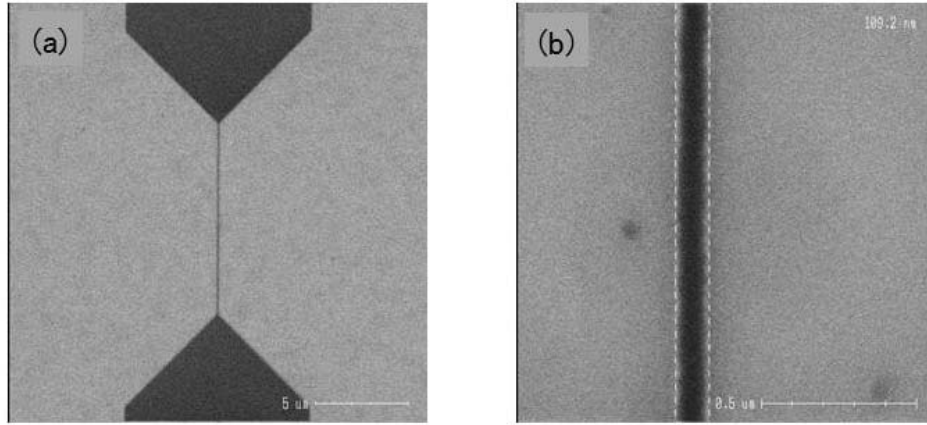


Figure 2. (a) SEM image of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ strip with the size of 100 nm wide \times 10 μm long \times 5 nm thick before milling and (b) an enlarged view.

3. I-V characteristics

Figure 3 shows the current-voltage (I-V) characteristics of the strip at various temperatures using two terminal methods. At 3 K, it shows a voltage jump at a critical current (I_c) and a small hysteresis behavior (inset). These behaviors are different to flux-flow type behaviors which are frequently observed in inhomogeneous cuprate nanostructures, and suggest the high uniformity of the strip. The I_c is 115 μA , which corresponds to the critical current density (J_c) of $2.3 \times 10^7 \text{ A/cm}^2$. The large J_c value also indicates the high quality of the strip. The I_c structure can be observed clearly up to 35 K, and the slope of I-V curve strongly decreases at 45 K corresponding to the normal resistance. These features clearly show that it is possible to fabricate the high quality $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ nanostructure with the size of 100 nm wide \times 10 μm long \times 5 nm thick without degradation.

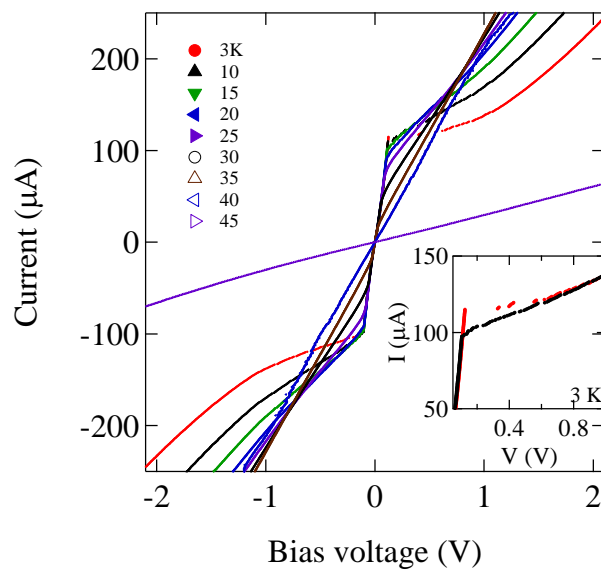


Figure 3. I-V characteristics of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ strip with the size of 100 nm wide \times 10 μm long \times 5 nm thick at various temperatures. Inset: Hysteresis behavior at 3 K.

4. Photoresponse

Electro-optical characterization was performed in a Gifford-McMahon cryocooler (SRDK-10HD, Sumitomo) between 3 K and 45 K. A bias current was supplied to the strip through a bias-T (PSPL5541A, Tektronix) without shunt resistor, and the signal was amplified by an rf amplifier (10KHz-1450MHz, NF2.0dB, LNA-1450, RF Bay) with a total gain of 30 dB and fed to a 20 GHz bandwidth single-shot oscilloscope (DPO72004, Textronix) [21]. The strip was illuminated with a femtosecond 1560 nm fiber laser (TC-1550, Menlo Systems) with a repetition rate of 100 MHz. The optical spot size was about 10 μm in diameter.

No signals are observed at zero bias current ($I_{\text{bias}} = 0$), but the signal appears as the bias current increases, and the height of the signal becomes maximum at the bias current just below I_c . Figure 4 shows the optical power dependence of the transient photoresponse of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ strip at 3 K. The I_{bias} is 111 μA , just below I_c . The height of the signal decreases as the optical intensity decreases, and disappears below -10 dBm. The estimated incident number of photons on the nanostrip is 9.9×10^4 photons/pulse at -10dBm, and it is far from single-photon detection regime.

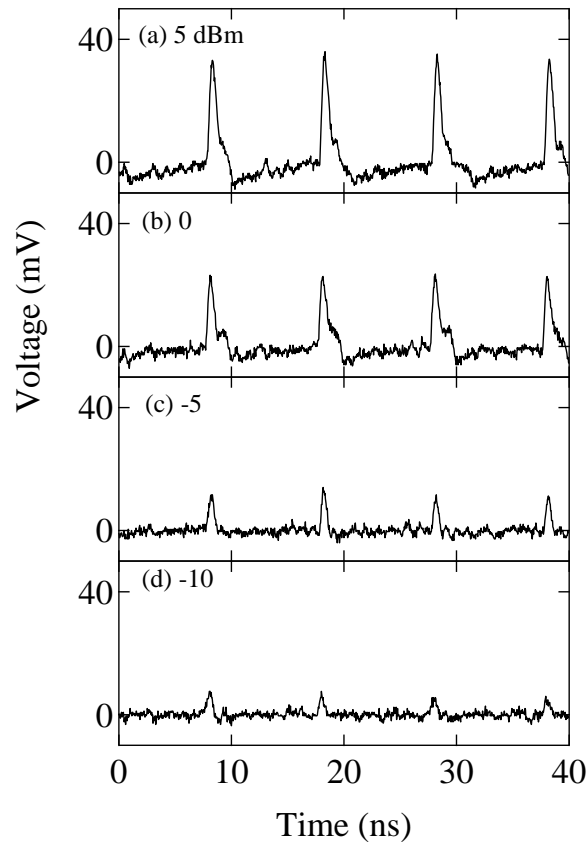


Figure 4. Transient photoresponse of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ strip with the size of 100 nm wide \times 10 μm long \times 5 nm thick with various laser pulse power. The strip is illuminated with a 1.56 μm wavelength femtosecond laser pulse at a 100 MHz repetition rate and measured at 3 K with $I_{\text{bias}} = 111 \mu\text{A}$.

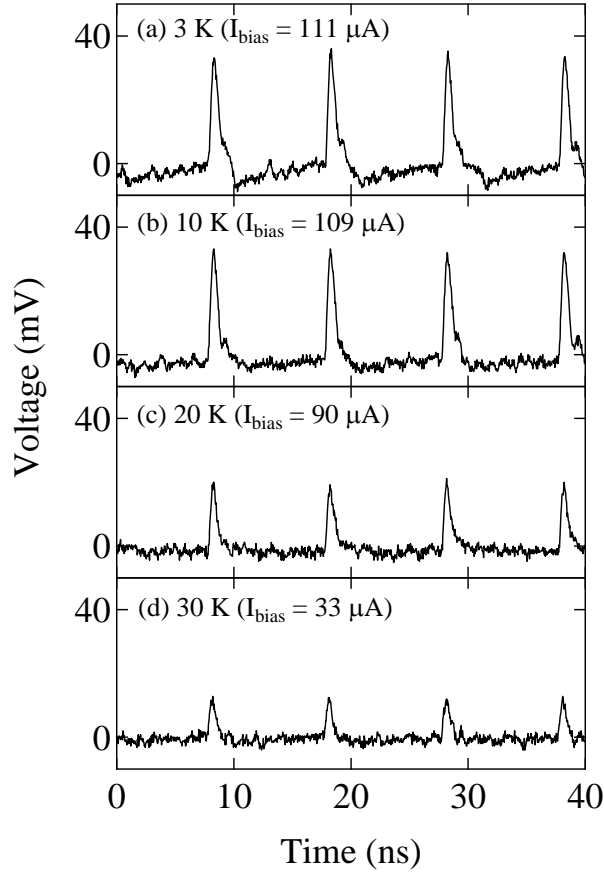


Figure 5. Transient photoresponse of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ strip with the size of $100 \text{ nm wide} \times 10 \text{ } \mu\text{m long} \times 5 \text{ nm thick}$ with various temperatures. The strip is illuminated with a $1.56 \text{ } \mu\text{m}$ wavelength femtosecond laser pulse at a 100 MHz repetition rate and I_{bias} 's are adjusted just below I_c .

Figure 5 shows the temperature dependence of the transient photoresponse of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ strip. Here, the bias currents are adjusted just below I_c at each temperature. As the temperature increases, the height of the signal decreases, but it can be observed up to 30 K . The decrease of the height is simply due to the decreases of I_{bias} .

Although the reason of these photoresponse is not clear, it seems to be a bolometric response of the nanostrip under the bias current. We notice that the small signals remain even when I_{bias} exceeds I_c , which may be due to the large temperature coefficient of the resistivity in the normal state. It is also noted that the same response have been also observed in the MgB_2 nanostrip with the width of 300 nm [22]. The response of MgB_2 nanostrip changes to the multi-photon detection regime with the width of 200 nm and changes to the single-photon detection regime with the width of 100 nm [8,22]. The bolometric response seems a common feature for superconducting nanostrip with the large cross-section.

Finally, we discuss the possibility towards single-photon detection of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ strip. Due to the high J_c of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$, the I_c of the present strip is about 5 times larger than the standard SSPD using other material (about 20 μA). So, we may have to reduce the cross section of our nanostrip about 5 times for single photon detection. The homogeneous nanostrip with 20 nm wide is required when the 5 nm thick $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ film is used. The requirement is quite challenging and needs further technological progresses. At the present stage, cuprate based SSPDs may be useful to detect high energy particles, such as single photon in the X ray region and single biomolecular ion [23,24]. Actually, in the case of MgB_2 , single biomolecular ion detection has been reported with 100 % detection efficiency up to 13 K [23].

5. Conclusion

We grow 5 nm thick $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ film using MBE and fabricate a strip with the size of 100 nm wide \times 5 nm thick with $J_c(3\text{K}) = 2.3 \times 10^7 \text{ A/cm}^2$. The I-V characteristics show a hysteresis and a sharp voltage jump at I_c . We observe the photoresponse of the strip under bias current even at 30 K, which is due to the bolometric response. We discuss the possible single-photon detection using $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$ nanostrip at high working temperature.

Acknowledgements

The authors would like to thank A. Tsukada, M. Naito, and H. Sato for many useful discussions about ultrathin film growth of $\text{La}_{1.85}\text{Sr}_{0.15}\text{CuO}_4$.

References

- [1] Natarajan C M, Tanner M G and Hadfield R H 2012 *Supercond. Sci. Technol.* **25**, 063001
- [2] Gol'tsman G N, Okunev O, Chulkova G, Lipatov A, Semenov A, Smirnov K, Voronov B, Dzardanov A, Williams C and Sobolewski R 2001 *Appl. Phys. Lett.* **79** 705–7
- [3] Marsili F, Verma V B, Stern J A, Harrington S, Lita A E, Gerrits T, Vayshenker I, Baek B, Shaw M D, Mirin R P and Nam S W 2013 *Nat. Photon.* **7**, 210
- [4] Zhang W J, You L X, Li H, Huang J, Lv C L, Zhang L, Liu X Y, Wu J J, Wang Z, Xie X M 2016 *arXiv*, 1609.00429
- [5] Shibata H, Shimizu K, Takesue H and Tokura Y 2013 *Appl. Phys. Express* **6**, 072801
- [6] Shibata H, Shimizu K, Takesue H and Tokura Y 2015 *Opt. Lett.* **40**, 3428
- [7] Shibata H, Honjo T and Shimizu K 2014 *Opt. Lett.* **39**, 5078
- [8] Shibata H, Takesue H, Honjo T, Akazaki T and Tokura Y 2010 *Appl. Phys. Lett.* **97** 212504
- [9] Shibata H 2014 *Appl. Phys. Express* **7**, 103101
- [10] Velasco A E, Cunnane D P, Acharya N, Briggs R, Beyer A, Shaw M, Karasik B S, Wolak M A, Xi X and Marsili F 2016 CLEO_QELS, FW4C.5
- [11] Curtz Z, Koller E, Zbinden H, Decroux M, Antognazza L, Fischer O and Gisin N 2010 *Supercond. Sci. Technol.* **23**, 045015
- [12] Atikian H A, Ghamsari B G, Anlage S M and Majedi A H 2011 *Appl. Phys. Lett.* **98** 081117
- [13] Arpaia R, Ejrnaes M, Parlato L, Cristiano R, Arzeo M, Bauch T, Nawaz S, Tafuri F, Pepe G P and Lombardi F 2014 *Supercond. Sci. Technol.* **27**, 044027
- [14] Lyatti M, Savenko A, Poppe U 2016 *Supercond. Sci. Technol.* **29**, 065017
- [15] Amari P, Palma C F, Jouan A, Couedo F, Bourlet N, Geron E, Malnou M, Mechin L, Sharafiev A, Lesueur J and Bergeal N 2016 *arXiv*. 1612.07730v2
- [16] Sato H, Yamamoto H and Naito N 1997 *Physica C* **274** 227
- [17] Sato H, Tsukada A, Naito M and Matsuda A 2000 *Phys. Rev.* **B61**, 12447
- [18] Logvenov G, Gozar A and Bozovic I 2009 *Science* **326**, 699
- [19] Shibata H, Karimoto S, Tsukada A and Makimoto T 2007 *J. Cryst. Growth* **301** 684
- [20] Shibata H, Akazaki T and Tokura Y 2013 *Appl. Phys. Express* **6**, 023101
- [21] Zhang L, Yan X, Jia X, Chen J, Kang L and Wu P 2017 *Appl. Phys. Lett.* **110**, 072602
- [22] Shibata H, Asahi M, Maruyama T, Akazaki T, Takesue H, Honjo T and Tokura Y 2009 *IEEE Trans. Appl. Supercond.* **19**, 358
- [23] Zen N, Shibata H, Mawatari Y, Koike M and Ohkubo M 2015 *Appl. Phys. Lett.* **106** 222601
- [24] Zhang X, Wang Q and Shilling A 2016 *AIP Advances* **6** 115104